

EQUULEUS Mission Analysis: Design of the Transfer Phase

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This paper highlights the design of the transfer phase of EQUULEUS (EQUilibriUm Lunar-Earth point 6U Spacecraft), which aims to reach and stay at the Earth-Moon L_2 libration point orbit to observe the far-side of the Moon. Since EQUULEUS is a piggyback CubeSat, there are original challenges that past missions have not faced so far. We develop a systematic trajectory design approach to tackle the difficulties. The trajectory is a nearly-ballistic low-energy transfer to a quasi-periodic halo orbit around Earth-Moon L_2 using lunar flybys and solar perturbation.

Key Words: EQUULEUS, Trajectory Design, Lunar Flyby, Solar Perturbation, Optimization

1. Introduction

Scheduled in 2018, Exploration Mission-1 (EM-1) will launch the Orion Multi-Purpose Crew Vehicle with 13 CubeSats via the NASA's new SLS rocket. EQUULEUS (EQUilibriUm Lunar-Earth point 6U Spacecraft, Fig. 1) is one of the selected CubeSats, proposed and being developed by JAXA and The University of Tokyo. EQUULEUS aims to reach a lunar libration point orbit around the Lagrange point L_2 to observe the far side of the Moon.^{1), 2), 3)}

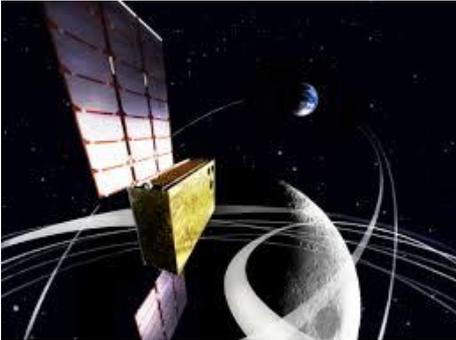


Fig. 1. An artist's concept of EQUULEUS.

In order to reach the lunar L_2 libration point orbit, EQUULEUS will use lunar flybys (LFBs) and solar perturbation to achieve low-energy transfer.⁴⁾ Other CubeSats in EM-1 such as Lunar IceCube⁵⁾, Lunar Polar Hydrogen Mapper⁶⁾, Lunar Flashlight⁷⁾, and OMOTENASHI⁸⁾ will transfer to the Moon, and Lunar Icecube and Lunar Polar Hydrogen Mapper also plan to use luni-solar perturbation.

In past missions, HITEN⁹⁾ and GRAIL¹⁰⁾ realized low-energy transfers to the Moon, and ARTEMIS¹¹⁾ reached the Earth-Moon L_2 libration point orbit by using LFBs and solar perturbation. Many studies investigated the class of

low-energy transfer in terms of the gravitational capture mechanism^{4), 12)}, the dynamical systems theory¹³⁾, and global search strategies^{14), 15), 16)}.

Comparing with these past missions and researches, EQUULEUS has its original challenges because it is a piggyback, CubeSat. The first challenge is the constrained launch condition of SLS, which targets low-altitude LFB making a ballistic trajectory escape from the Earth-Moon system. Therefore, EQUULEUS must change the 1st LFB condition by its thruster to return to the Earth-Moon system. The demonstration of the controllability of a CubeSat under luni-solar perturbation is one of the main objectives. A robust and flexible framework of designing transfer trajectories is necessary to deal with constrained initial conditions especially near launch.

The second challenge is the limited Δv budget, approximately 80 [m/s] for the whole mission, including trajectory correction maneuvers (TCMs) and clean up maneuvers (CUMs) before and after LFBs, and stationkeeping Δv on the science orbit²¹⁾. Ideally, the transfer phase is nearly ballistic for arbitrary launch conditions.

The third challenge is the critical execution of Δv_1 to achieve the desirable 1st LFB condition. Currently, EQUULEUS plans to start Δv_1 1 day after the separation from SLS and finish the execution of Δv_1 , orbit determination, and TCMs before the 1st LFB, which occurs approximately 6 days after the separation. Considering the limited thrust capability, 3 [mN] at most, the operation before the 1st FLB could be tight²⁾ and therefore it is important to reduce Δv_1 as much as possible.

This paper presents our current strategy to tackle these challenges. We develop a systematic method for designing transfer trajectories from the separation from SLS to capture into a lunar L_2 quasi-periodic halo orbit in a high-fidelity

model. Error analyses including TCM, CUM, and Monte Carlo simulation are out of scope of this paper. The present analysis uses the launch condition that the epoch of the separation is 2018 OCT 07 15:39 (UTC) given by NASA.

The remainder of this paper is organized as follows. Section 2 introduces the basic concept of low-energy transfer. Section 3 highlights our approach of designing transfer trajectories. Section 4 shows results of computing transfer trajectories.

2. Low-Energy Transfer

Belbruno and Miller⁴⁾ found low-energy transfer to the Moon, which has remarkable benefits such as lower Δv and a wider launch window as compared with the high-energy Hohmann-type transfer. The mechanism of the low-energy lunar transfer is based on the use of the solar tidal force¹²⁾, which accelerates a trajectory, pumps up its perigee, and reduces v_∞ with respect to the Moon if an apogee is placed in the 2nd or 4th quadrant around the Earth and vice versa.

Fig. 2 shows an example of low-energy transfer to the Moon¹⁶⁾, which places its apogee in the 2nd quadrant to reduce v_∞ . The trajectory uses LFB to save Earth departure Δv and exploits solar perturbation to reduce lunar insertion Δv .

Since the original trajectory of Orion encounters the Moon with high v_∞ , EQUULEUS needs to reduce v_∞ by using solar perturbation to achieve low-energy capture into the Earth-Moon L_2 quasi-periodic halo orbit.

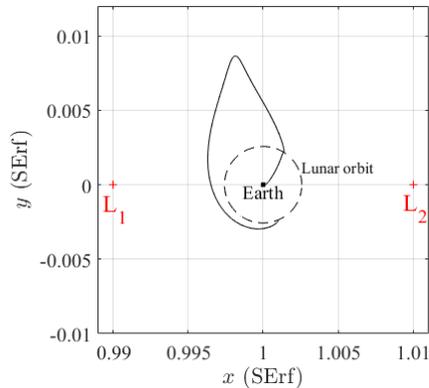


Fig. 2. An example of low-energy transfer to the Moon computed in Oshima et al.¹⁶⁾ The trajectory is shown in the non-dimensional Sun-Earth rotating frame.

3. Trajectory Design Approach

Fig. 3 shows a flowchart of our trajectory design approach for EQUULEUS. The following sections highlight each process with some backgrounds.

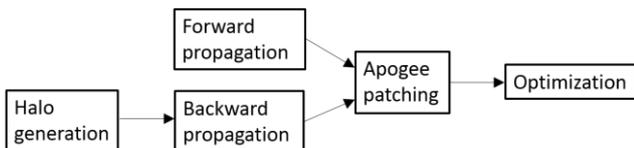


Fig. 3. A flowchart of the trajectory design approach for EQUULEUS.

3.1. Halo Generation Process

This process computes multi-revolutional quasi-periodic orbits around lunar L_2 . Periodic halo orbits in the Earth-Moon circular restricted three-body problem (CR3BP) are optimized along the wide range of epoch in the full-ephemeris model by using the optimization software jTOP.¹⁷⁾

Note that periodic orbits in autonomous systems such as the CR3BP cannot persist when including time-dependent perturbations except for special resonant cases. Instead, natural dynamical substitutes of periodic orbits in non-autonomous systems such as the ephemeris model are quasi-periodic orbits.^{18), 19), 20)}

See Oguri et al.²¹⁾ for details of the computation of quasi-periodic halo orbits and the analysis of stationkeeping Δv for EQUULEUS. Fig. 4 shows eight families of quasi-periodic halo orbits computed in Oguri et al.²¹⁾ In this paper, we use the quasi-periodic halo orbits of the epoch from 2019 JAN 01 to 2019 MAY 31, but those of a different time range could generate more families of solutions.

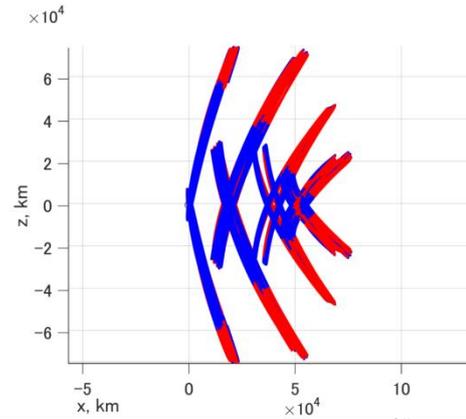


Fig. 4. Families of quasi-periodic halo orbits.²¹⁾

3.2. Forward Propagation Process

This process propagates trajectories forward in time from the separation, currently including gravitational effects of the Sun, Earth, and Moon as point masses, and saves time and states at apogees. After 1 day from the separation, we apply impulsive Δv_1 to change the 1st LFB condition. The computation is three-dimensional grid search in terms of the magnitude and angles of Δv_1 . Since the execution of Δv_1 is critical as denoted in Section 1, we set the lower bound of the magnitude of Δv_1 such that some trajectories can come back to the Earth-Moon system after the 1st apogee even though most of them escape. In this analysis, the lower bound is 2 [m/s] and the upper bound is 6 [m/s].

Fig. 5 shows an example of a trajectory propagated forward in time with $\Delta v_1=2$ [m/s]. Since Δv_1 is small, the trajectory comes back to the Earth-Moon system after the high 1st apogee distance. We confirm that the large values of z ($> 770,000$ [km]) at 1st apogees of $\Delta v_1=2$ [m/s] make it difficult to patch with backward legs from the quasi-periodic halo orbits (see Section 3.3), i.e., larger Δv_1 is necessary to design trajectories for the conditions of the current analysis.

Fig. 6 shows values of (a) x , (b) y , and (c) z at apogees in the Earth-centered J2000 frame in terms of ephemeris time as the result of the grid search. From these figures, we extract

time and position of the 1st apogee as

$$\begin{aligned} 5.95 \times 10^8 \leq t \leq 5.995 \times 10^8 \text{ [s]}, \\ -12.5 \times 10^5 \leq x \leq -5 \times 10^5 \text{ [km]}, \\ -3.8 \times 10^5 \leq y \leq 16 \times 10^5 \text{ [km]}, \\ 0 \leq z \leq 12 \times 10^5 \text{ [km]}. \end{aligned} \quad (1)$$

This condition of the 1st apogee is used to restrict the data to patch at 1st apogees (see Section 3.4).

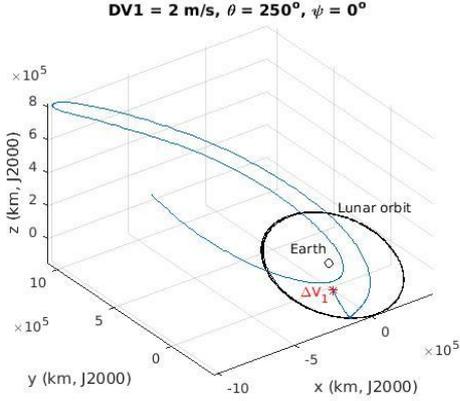


Fig. 5. An example of a trajectory propagated forward in time with $\Delta v_1=2$ [m/s] in the Earth-centered J2000 frame.

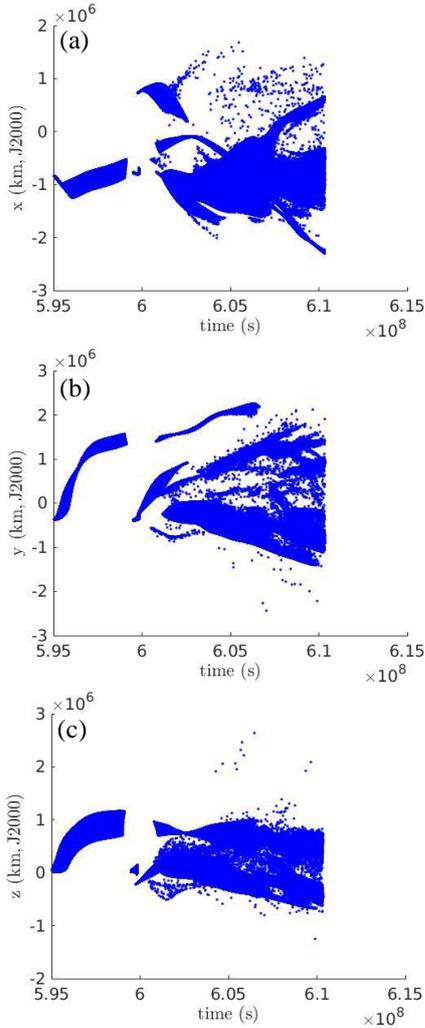


Fig. 6. (a) x, (b) y, (c) z at apogees of trajectories propagated forward in time in the Earth-centered J2000 frame in terms of ephemeris time.

3.3. Backward Propagation Process

This process propagates trajectories backward in time from quasi-periodic halo orbits, currently including gravitational effects of the Sun, Earth, and Moon as point masses, and saves time and states at apogees. In order to escape from the quasi-periodic halo orbits efficiently, we apply small Δv to the minimum stretching direction.

The minimum stretching direction δx_i is defined by using the state transition matrix (STM) Φ as

$$\min_{\delta x_i} \|\delta x_f\|, \quad (2)$$

where

$$\|\delta x_f\| = \sqrt{\delta x_i^T \Phi^T \Phi \delta x_i}. \quad (3)$$

From Eqs. (2) and (3), the minimum stretching direction is the eigenvector associated to the minimum eigenvalue of the Cauchy-Green tensor $C = \Phi^T \Phi$.

The minimum stretching direction of the Cauchy-Green tensor can be regarded as an extension of the eigenvector of the stable direction of the Monodromy matrix in autonomous systems to non-autonomous systems, which was used in the computation of Lagrangian coherent structures^{22), 23)} and an analysis of the low-energy gravitational capture of BepiColombo to Mercury.²⁴⁾

Fig. 7 shows trajectories propagated backward in time by applying small Δv ($=0.1$ [m/s]) to minimum stretching directions from one quasi-periodic halo orbit. The trajectories can be regarded as an extension of stable manifolds in autonomous systems to non-autonomous systems. EQUULEUS aims to get on one of these time-dependent analogues of stable manifolds²³⁾ to be captured into a quasi-periodic halo orbit with small insertion Δv .

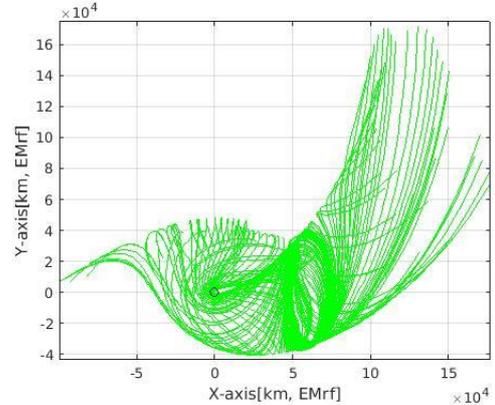


Fig. 7. Trajectories propagated backward in time by applying small velocity perturbation to minimum stretching directions from a quasi-periodic halo orbit around lunar L₂ in the Moon-centered Earth-Moon rotating frame.

3.4. Apogee Patching Process

This process patches forward (Section 3.2) and backward (Section 3.3) legs at 1st apogees. We only use the data at apogees of the backward propagation that satisfy Eq. (1). We use the MATLAB function `knnsearch` to find k-nearest neighbors in terms of time and position at apogees, and extract pairs such that differences of time and distance between apogees are less than prescribed tolerances (TOLE.t and TOLE.rmag). We sort these pairs in ascending order in terms of the magnitude of the difference of velocity (discontinuity

Δv). If the differences of time and distance between apogees are small enough, discontinuity Δv could be a good criterion to measure the quality of initial guesses, which is ideally equal to Δv subtracting Δv_1 from total Δv after optimization.

Fig. 8 shows forward ($\Delta v_1=6$ [m/s], blue) and backward (red) states at apogees, displayed with those satisfying tolerances $\text{TOLE.t}=0.025$ [days] and $\text{TOLE.rmag}=25,000$ [km] (green for forward and black for backward) in (a) x-y, (b) x-z, (c) vx-vy, (d) vx-vz planes in the Earth-centered J2000 frame.

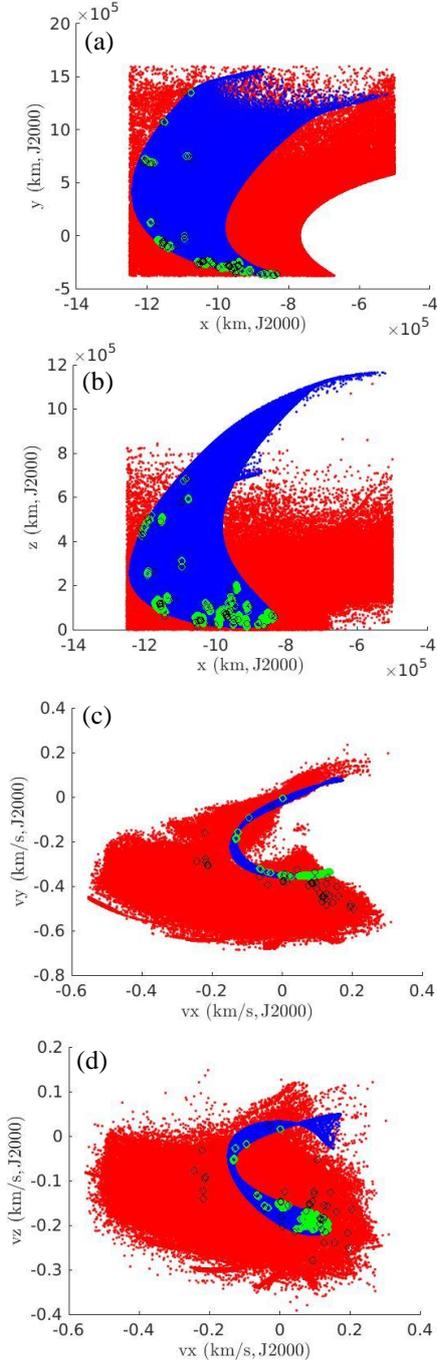


Fig. 8. Forward ($\Delta v_1=6$ [m/s], blue) and backward (red) states at apogees, displayed with those satisfying the tolerances $\text{TOLE.t}=0.025$ [days] and $\text{TOLE.rmag}=25,000$ [km] (green for forward and black for backward) in (a) x-y, (b) x-z, (c) vx-vy, (d) vx-vz planes in the Earth-centered J2000 frame.

We note that our previous method patched forward and backward legs at perilunes.^{2), 25), 26), 27)} However, the current method of patching at apogees has the following advantages as compared with the previous method. Firstly, the current method enables easier design of transfer trajectories using LFB once, i.e., only the 1st LFB, because the states at 1st LFBs of forward propagation are similar and it is difficult to distinguish good initial guesses. Note that substantial differences arise after the 1st LFB due to the sensitivity of LFB. Greater numbers of LFBs result in larger navigation Δv such as TCM and CUM, and more challenging operations. Secondly, the current method can naturally include deep space maneuver (DSM) as the velocity discontinuity at apogees in initial guess, whereas the previous method cannot without fixing Δv_1 or adding extra dimensionality to the search space. Thirdly, it is easier to predict Δv to patch forward and backward legs before optimization in the current method (see Section 4), whereas it was difficult in the previous method because states at perilunes are sensitive due to the strong gravity of the Moon^{26), 27)}.

3.5. Optimization Process

This module optimizes the initial guesses satisfying tolerances TOLE.t and TOLE.rmag by using the optimization software jTOP¹⁷⁾, which was successfully used to design the trajectory of PROCYON²⁸⁾, the first deep-space micro-spacecraft. jTOP covers from low-fidelity to high-fidelity dynamics and both impulsive and finite thrust maneuvers. Currently, we optimize trajectories of EQUULEUS including gravitational effects of the Sun, Jupiter, Earth, and Moon as point masses, and minimize total Δv by assuming impulsive maneuvers.

jTOP optimizes trajectories via the direct multiple shooting method. Fig. 9 schematically shows the patching strategy of jTOP, where arcs are propagated forward and backward in time from each phase, and additional phase is introduced in the sensitive region to improve the convergence.

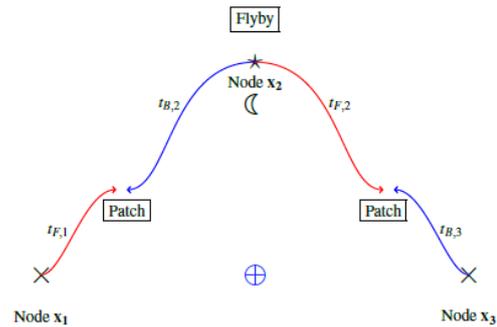


Fig. 9. A Schematic figure of the patching strategy of jTOP.

4. Results

This section presents the results of designing transfer trajectories for EQUULEUS based on the above procedures.

Fig. 10 shows the results of optimizing initial guesses of discontinuity $\Delta v < 100$ [m/s]. In this analysis, we only optimize one initial guess in each family, and families are classified in terms of dates of LFBs and time of flight (TOF). Fig. 10 (a) shows the values of TOF and total Δv . Several

families of solutions and trade-off between TOF and total Δv are visible. Fig. 10 (b) shows Δv_1 and total Δv , which indicates that the changes of Δv_1 via the optimization are small. It is possible to put a constraint on the magnitude of Δv_1 though we did not use it in this analysis. Fig. 10 (c) shows the discontinuity Δv at apogees of initial guesses and Δv subtracting Δv_1 from total Δv after optimization. There is a near-linear relationship between these values and thus Δv is quantitatively predictable before the optimization.

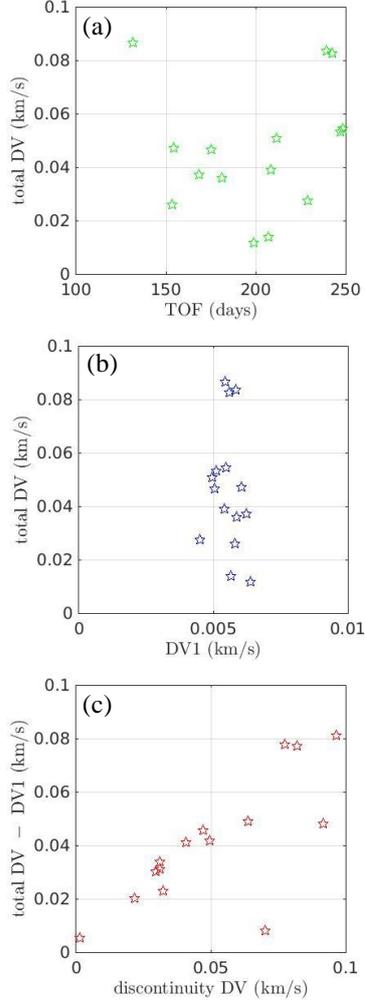


Fig. 10. (a) TOF and total Δv of optimal solutions. (b) Δv_1 and total Δv of optimal solutions. (c) Discontinuity Δv of the initial guesses and Δv subtracting Δv_1 from total Δv after optimization.

Fig. 11 (a), (b) show the transfer trajectory of the optimal solution with total $\Delta v=11.5$ [m/s], $\Delta v_1=6.4$ [m/s], TOF=199 [days] in the (a) Sun-Earth and (b) Earth-Moon rotating frames. The red and blue arcs represent the initial guess, and the green and magenta arcs represent the optimal solution. Since the discontinuity at the apogee of the initial guess was small, the trajectories of the initial guess and the optimal solution nearly overlap. The trajectory reaches a quasi-periodic halo orbit by using LFB only once, which could save navigation Δv and relax complexity of the operation.

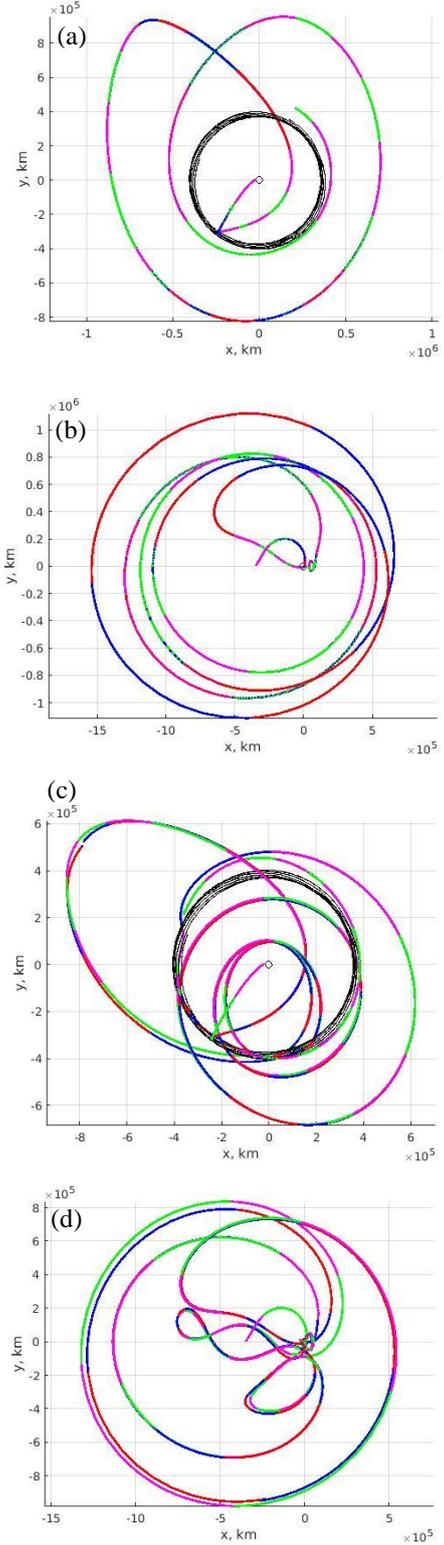


Fig. 11. Optimal transfer trajectories with (a), (b) total $\Delta v=11.5$ [m/s], $\Delta v_1=6.4$ [m/s], TOF=199 [days], and (c), (d) total $\Delta v=13.7$ [m/s], $\Delta v_1=5.7$ [m/s], TOF=207 [days] in the (a), (c) Sun-Earth and (b), (d) Earth-Moon rotating frames, respectively. The red and blue arcs represent the initial guess, and the green and magenta arcs represent the optimal solution.

Fig. 11 (c), (d) show the transfer trajectory of the optimal solution with total $\Delta v=13.7$ [m/s], $\Delta v_1=5.7$ [m/s], TOF=207 [days] in the (c) Sun-Earth and (d) Earth-Moon rotating frames. This solution uses multiple LFBs, but results in smaller Δv_1 . Multiple LFBs may be useful to reduce Δv_1 , and there could be trade-off between numbers of LFBs and the magnitude of Δv_1 .

5. Conclusion

This paper presented the trajectory design approach of EQUULEUS for low-energy transfers to the Earth-Moon L_2 quasi-periodic halo orbit. Our method could systematically compute families of transfers using lunar flybys and solar perturbation. We showed that total Δv and Δv_1 are predictable before optimization, which is a favorable feature for an efficient search.

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